

ISSN 2518-1726 (Online),  
ISSN 1991-346X (Print)

ҚАЗАҚСТАН РЕСПУБЛИКАСЫ  
ҰЛТТЫҚ ҒЫЛЫМ АКАДЕМИЯСЫНЫҢ  
Әль-фараби атындағы Қазақ ұлттық университетінің

# Х А Б А Р Л А Р Ы

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## ИЗВЕСТИЯ

НАЦИОНАЛЬНОЙ АКАДЕМИИ НАУК  
РЕСПУБЛИКИ КАЗАХСТАН  
Қазақстан Республикасының  
Ғылым Академиясының  
Әль-Фараби атындағы  
Қазақ ұлттық университетінің

## NEWS

OF THE NATIONAL ACADEMY OF SCIENCES  
OF THE REPUBLIC OF KAZAKHSTAN  
Al-farabi kazakh  
national university

**SERIES**  
**PHYSICO-MATHEMATICAL**

**2 (324)**

**MARCH - APRIL 2019**

PUBLISHED SINCE JANUARY 1963

PUBLISHED 6 TIMES A YEAR

ALMATY, NAS RK

Б а с р е д а к т о р ы  
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ISSN 2518-1726 (Online), ISSN 1991-346X (Print)

Меншіктенуші: «Қазақстан Республикасының Ұлттық ғылым академиясы» РҚБ (Алматы қ.)  
Қазақстан республикасының Мәдениет пен ақпарат министрлігінің Ақпарат және мұрағат комитетінде  
01.06.2006 ж. берілген №5543-Ж мерзімдік басылым тіркеуіне қойылу туралы куәлік

Мерзімділігі: жылына 6 рет.  
Тиражы: 300 дана.

Редакцияның мекенжайы: 050010, Алматы қ., Шевченко көш., 28, 219 бөл., 220, тел.: 272-13-19, 272-13-18,  
<http://physics-mathematics.kz/index.php/en/archive>

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Типографияның мекенжайы: «Аруна» ЖК, Алматы қ., Муратбаева көш., 75.

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«Известия НАН РК. Серия физико-математическая».

ISSN 2518-1726 (Online), ISSN 1991-346X (Print)

Собственник: РОО «Национальная академия наук Республики Казахстан» (г. Алматы)

Свидетельство о постановке на учет периодического печатного издания в Комитете информации и архивов  
Министерства культуры и информации Республики Казахстан №5543-Ж, выданное 01.06.2006 г.

Периодичность: 6 раз в год.

Тираж: 300 экземпляров.

Адрес редакции: 050010, г. Алматы, ул. Шевченко, 28, ком. 219, 220, тел.: 272-13-19, 272-13-18,  
<http://physics-mathematics.kz/index.php/en/archive>

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**News of the National Academy of Sciences of the Republic of Kazakhstan. Physical-mathematical series.**

**ISSN 2518-1726 (Online), ISSN 1991-346X (Print)**

Owner: RPA "National Academy of Sciences of the Republic of Kazakhstan" (Almaty)

The certificate of registration of a periodic printed publication in the Committee of information and archives of the Ministry of culture and information of the Republic of Kazakhstan N 5543-Ж, issued 01.06.2006

Periodicity: 6 times a year

Circulation: 300 copies

Editorial address: 28, Shevchenko str., of. 219, 220, Almaty, 050010, tel. 272-13-19, 272-13-18,  
<http://physics-mathematics.kz/index.php/en/archive>

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Address of printing house: ST "Aruna", 75, Muratbayev str, Almaty

## NEWS

OF THE NATIONAL ACADEMY OF SCIENCES OF THE REPUBLIC OF KAZAKHSTAN  
PHYSICO-MATHEMATICAL SERIES

ISSN 1991-346X

<https://doi.org/10.32014/2019.2518-1726.14>

Volume 2, Number 324 (2019), 69 – 78

UDC 531, 532.133, 621.3.018.72.025.1

IRSTI 29.03.77

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## CALCULATION AND VISUALIZATION OF SMALL OSCILLATIONS OF A DOUBLE PLANE PENDULUM

**Abstract.** The article considers the calculation and visualization of small oscillations of a double plane simple pendulum. It contains the brief derivation of the motion equation and its solutions; the mathematical model of the motion in the form of the system of nonlinear differential equations. Experiments with the double pendulum are performed at various ratios of bodies' masses and initial angles, in particular for three values of  $\mu = m_1 / m_2$ :  $\mu_1 = 0.1$ ,  $\mu_2 = 0.2$ ,  $\mu_3 = 0.3$  when the lengths of the pendulum are  $l = l_1 = l_2 = 0.25$  m and  $g = 9.8$  m/s<sup>2</sup>. The angles of pendulum deviation are given in radians. The graphs of small oscillations of the pendulum show that beats occur in the system during which energy cyclically passes from one pendulum to another. When one pendulum almost stops, the other pendulum is at its maximum amplitude. After a while pendulums "exchange their states" and so on. Oscillations with a bigger frequency  $\omega_1$  are modulated by lower frequency oscillations with a frequency  $\omega_2$ .

**Key words.** Double pendulum, small oscillations, beats, energy exchange, natural frequencies, natural mode.

**Introduction.** Nowadays all educational institutions of Kazakhstan are provided with computer hardware and software, interactive boards and internet. Almost all teachers have completed language and computer courses for professional development. Hence the educational institutions have all conditions for using computer training programs and models for performing computer laboratory works. In recent years the new computer system for carrying out mathematical calculations MATLAB is being widely used in many universities and engineering institutions throughout the world [1-7]. Unfortunately, the numerical calculations carried out by students are often done by means of the calculator. Modern computers are frequently used only for presentation of the work. Actually students should be able not only to solve these or other engineering problems, but also do them by using modern methods, that is, using personal computers.

Students of the physics specialties 5B060400 and 5B011000 successfully master the discipline "Computer modeling of physical phenomena" which is the logical continuation of the disciplines "Information technologies in teaching physics" and "Use of electronic textbooks in teaching physics". The goal of this discipline is to study and learn the MATLAB program language, acquaintance with its huge opportunities for modeling and visualization of physical processes.

In our early works [8-26] we used the MATLAB system for modeling and visualization of physical processes related with mechanics, molecular physics, electromagnetism and quantum physics. This software has enabled us to solve ordinary differential equations (ODE), visualize equipotential lines of charged conductors system, describe the motion of charged particles in electric, magnetic and gravitational fields and etc.

The present article is devoted to simulation of the double plane pendulum by using the package of MATLAB applied programs.

Double pendulum is, undoubtedly, a real miracle of the nature. The jump of complexity which is observed upon transition from a simple pendulum to a double one is amazing. Oscillations of the simple pendulum are of periodic nature. At small deviations from equilibrium such oscillations are harmonic and obey the sine or cosine law. In case of nonlinear oscillations the period depends on amplitude, but the periodicity of motion is conserved. In other words, the approximation of small oscillations quite perfectly reflects the essential properties of the simple pendulum.

The double pendulum "behaves" itself absolutely differently. At the mode of small oscillations the double pendulum moves in a beat pattern, i.e. with variation of resultant amplitude with time. At the increase of energy the nature of oscillations of the pendulum changes dramatically – oscillations become chaotic. In spite of the fact that the double pendulum can be described by the system of several ordinary differential equations, which is quite deterministic model, the emergence of chaos seems very unusual. This situation reminds the Lorentz system where the deterministic model of three equations also reveals chaotic behavior. Let's perform an experiment with the appendix given below and observe the motion of the double pendulum at various ratios of bodies' masses and initial angles.

At first we will develop the mathematical model of the double pendulum in the form of the system of nonlinear differential equations. Let us begin with the derivation of Lagrange equations.

**Lagrange equations.** In Lagrangian mechanics the system is described by using the generalized coordinates and the generalized velocities. In the considered problem the angles of deviation of the pendulums  $\alpha_1, \alpha_2$  and angular velocities  $\dot{\alpha}_1, \dot{\alpha}_2$  are taken as such variables. Using the specified variables, let us to derive the Lagrangian of the double pendulum and write down differential equations of Lagrange.

The simplified model of the double pendulum is shown in figure 1. Let's consider rods to be weightless. Their lengths are equal to  $l_1$  and  $l_2$ . The masses of small bodies (they are presented by balls of finite radius) are  $m_1$  and  $m_2$ . It is assumed that there is no friction in suspension points.

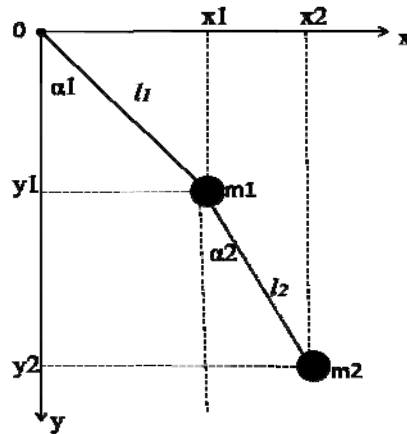


Fig.1 - The simplified model of the double plane pendulum

The masses' coordinates:

$$\begin{aligned} x_1 &= l_1 \sin \alpha_1, & x_2 &= l_1 \sin \alpha_1 + l_2 \sin \alpha_2, \\ y_1 &= -l_1 \cos \alpha_1, & y_2 &= -l_1 \cos \alpha_1 - l_2 \cos \alpha_2 \end{aligned} \quad (1)$$

The Lagrange function which is equal to the difference between kinetic and potential energies of the pendulum (respectively T and V) is expressed by the formula:

$$\begin{aligned}
 L = T - V &= \frac{m_1 v_1^2}{2} + \frac{m_2 v_2^2}{2} - V_1 - V_2 = \\
 &= \frac{m_1(\dot{x}_1^2 + \dot{y}_1^2)}{2} + \frac{m_2(\dot{x}_2^2 + \dot{y}_2^2)}{2} - m_1 g y_1 - m_2 g y_2
 \end{aligned} \tag{2}$$

By taking into account the formula (1) we receive the following expression of the Lagrange function:

$$\begin{aligned}
 L &= \left( \frac{m_1}{2} + \frac{m_2}{2} \right) l_1^2 \dot{\alpha}_1^2 + \frac{m_2}{2} l_2^2 \dot{\alpha}_2^2 + m_2 l_1 l_2 \dot{\alpha}_1 \dot{\alpha}_2 \cos(\alpha_1 - \alpha_2) + \\
 &+ m_1 g l_1 \cos \alpha_1 + m_2 g l_1 \cos \alpha_1 + m_2 g l_2 \cos \alpha_2
 \end{aligned} \tag{3}$$

Since we consider small oscillations the angles  $\alpha_1, \alpha_2$  are small so the  $\cos(\alpha_1 - \alpha_2)$  can be expanded into Maclaurin series:

$$\cos(\alpha_1 - \alpha_2) = 1 - \frac{(\alpha_1 - \alpha_2)^2}{2!} + \frac{(\alpha_1 - \alpha_2)^4}{4!} - \dots$$

We approximate the function  $\cos(\alpha_1 - \alpha_2)$  by taking the first term of a Maclaurin series since the subsequent terms are negligibly small. With this approximation the equation (3) becomes

$$\begin{aligned}
 L &= \left( \frac{m_1}{2} + \frac{m_2}{2} \right) l_1^2 \dot{\alpha}_1^2 + \frac{m_2}{2} l_2^2 \dot{\alpha}_2^2 + m_2 l_1 l_2 \dot{\alpha}_1 \dot{\alpha}_2 + \\
 &+ m_1 g l_1 \cos \alpha_1 + m_2 g l_1 \cos \alpha_1 + m_2 g l_2 \cos \alpha_2
 \end{aligned} \tag{4}$$

Now we can write down the Lagrange equation:

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\alpha}_i} \right) - \frac{\partial L}{\partial \alpha_i} = 0; \quad i = 1, 2 \tag{5}$$

Do the derivatives:

$$\frac{\partial L}{\partial \dot{\alpha}_1} = (m_1 + m_2) l_1^2 \dot{\alpha}_1 + m_2 l_1 l_2 \dot{\alpha}_2;$$

$$\frac{\partial L}{\partial \dot{\alpha}_2} = m_2 l_2^2 \dot{\alpha}_2 + m_2 l_1 l_2 \dot{\alpha}_1;$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\alpha}_1} \right) = (m_1 + m_2) l_1^2 \ddot{\alpha}_1 + m_2 l_1 l_2 \ddot{\alpha}_2;$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\alpha}_2} \right) = m_2 l_2^2 \ddot{\alpha}_2 + m_2 l_1 l_2 \ddot{\alpha}_1;$$

$$\frac{\partial L}{\partial \alpha_1} = -m_1 g l_1 \sin \alpha_1 - m_2 g l_1 \sin \alpha_1 = |\sin \alpha \approx \alpha| = -m_1 g l_1 \alpha_1 - m_2 g l_1 \alpha_1;$$

$$\frac{\partial L}{\partial \alpha_2} = -m_2 g l_2 \sin \alpha_2 = |\sin \alpha \approx \alpha| = -m_2 g l_2 \alpha_2;$$

Put it all together to get two Lagrange's differential equations:

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\alpha}_1} \right) - \frac{\partial L}{\partial \alpha_1} = (m_1 + m_2) l_1^2 \ddot{\alpha}_1 + m_2 l_1 l_2 \ddot{\alpha}_2 + (m_1 + m_2) g l_1 \alpha_1 = 0$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\alpha}_2} \right) - \frac{\partial L}{\partial \alpha_2} = m_2 l_2^2 \ddot{\alpha}_2 + m_2 l_1 l_2 \ddot{\alpha}_1 + m_2 g l_2 \alpha_2 = 0$$

This system of differential equations can be written in a compact matrix form. Firstly, we introduce the following matrixes:

$$\alpha(t) = \begin{pmatrix} \alpha_1(t) \\ \alpha_2(t) \end{pmatrix}, \quad M = \begin{pmatrix} (m_1 + m_2) l_1^2 & m_2 l_1 l_2 \\ m_2 l_1 l_2 & l_2^2 m_2 \end{pmatrix},$$

$$K = \begin{pmatrix} (m_1 + m_2) g l_1 & 0 \\ 0 & m_2 g l_2 \end{pmatrix}, \quad 0 = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

Then the system of differential equations is expressed as

$$M \ddot{\alpha} + K \alpha = 0$$

When there is the oscillation of one body such equation describes free undamped oscillation with a particular frequency. In case of the double pendulum oscillation the solution of this equation (as you will see below) will contain two characteristic frequencies which are called normal or natural modes. Natural modes represent the real part of a complex vector function

$$\alpha(t) = \begin{pmatrix} \alpha_1(t) \\ \alpha_2(t) \end{pmatrix} = \text{Re} \begin{bmatrix} (H_1) e^{i\omega t} \\ (H_2) e^{i\omega t} \end{bmatrix}$$

here  $H_1, H_2$  are the eigenvectors,  $\omega$  is the real value of the frequency. The values of natural frequencies  $\omega_{1,2}$  are determined by solving the characteristic equations

$$\det(K - \omega^2 M) = 0$$

$$\begin{bmatrix} (m_1 + m_2) g l_1 - \omega^2 (m_1 + m_2) l_1^2 & -\omega^2 m_2 l_1 l_2 \\ -\omega^2 m_2 l_1 l_2 & m_2 g l_2 - \omega^2 m_2 l_2^2 \end{bmatrix}$$

$$(m_1 + m_2) g^2 - \omega^2 (m_1 + m_2) (l_1 + l_2) g + \omega^4 m_1 l_1 l_2 = 0$$

We obtained the biquadratic equation for frequencies  $\omega$ . Let us calculate the discriminant:

$$\begin{aligned} D &= (m_1 + m_2)^2 g^2 (l_1 + l_2)^2 - 4 m_1 (m_1 + m_2) g^2 l_1 l_2 = \\ &= g^2 (m_1 + m_2) \left[ (m_1 + m_2) (l_1 + l_2)^2 - 4 m_1 l_1 l_2 \right]; \end{aligned}$$

Thus the square of the natural frequencies  $\omega_{1,2}$  is equal to

$$\omega_{1,2}^2 = \frac{g}{2 m_1 l_1 l_2} \left\{ (m_1 + m_2) (l_1 + l_2) \pm \sqrt{(m_1 + m_2) \left[ (m_1 + m_2) (l_1 + l_2)^2 - 4 m_1 l_1 l_2 \right]} \right\}$$



This expression is quite lengthy. Therefore we assume that the lengths of both pendulums are the same:  $l_1 = l_2 = l$ . Then the natural frequencies will be determined by more compact formula

$$\omega_{1,2}^2 = \frac{g}{l} \left[ (1 + \mu) \pm \sqrt{(1 + \mu)\mu} \right], \text{ where } \mu = \frac{m_2}{m_1}$$

This formula shows that natural frequencies  $\omega_{1,2}$  depend only upon the parameter  $\mu$  (when  $g/l = 1$ ). If masses are identical  $m_1 = m_2 = m$ , i.e.  $\mu = 1$  then the natural frequencies are equal to

$$\omega_{1,2} = \sqrt{\frac{g}{l}(2 \pm \sqrt{2})}.$$

Here is the MatLab program for calculation and visualization of the double plane pendulum's oscillation:

```
>> mye=0.1;
>> g=9.8;
>> l=0.25;
>> A=mye+((1+mye)*mye).^0.5;
>> w1=(g/l).^0.5*(1+A).^0.5;
>> w2=(g/l).^0.5*(1-A).^0.5;
>> t=0:0.1:5;
>> alfa1=(-pi/12)*(mye./(1+mye)).^0.5*(cos(w1.*t))+... (pi/12)*(mye./(1+mye)).^0.5*(cos(w2.*t));
>> plot(t,alfa1)
>> grid on
>> hold on
>> alfa2=(pi/12)*cos(w1.*t)+(pi/12)*cos(w2.*t);
>> plot(t,alfa2,'ro-')
>> gtext('alfa2')
>> gtext('alfa1')
>> gtext('mye=0.1')
```

The result is presented in the fig.2

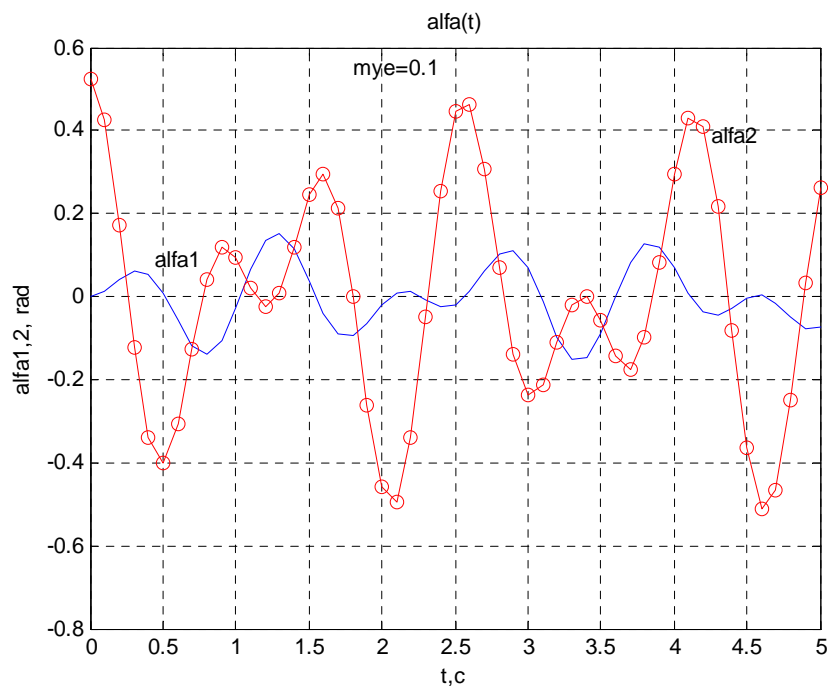


Fig.2 - The graphs of small oscillations of the double pendulum at  $\mu_1 = 0.1$

```

>> mye=0.2;
>> g=9.8;
>> l=0.25;
>> A=mye+((1+mye)*mye).^0.5;
>> w1=(g/l).^0.5*(1+A).^0.5;
>> w2=(g/l).^0.5*(1-A).^0.5;
>> t=0:0.1:5;
>> alfa1=(-pi/12)*(mye./(1+mye)).^0.5*(cos(w1.*t))+...
(pi/12)*(mye./(1+mye)).^0.5*(cos(w2.*t));
>> plot(t,alfa1)
>> grid on
>> hold on
>> alfa2=(pi/12)*cos(w1.*t)+(pi/12)*cos(w2.*t);
>> plot(t,alfa2,'ro-')
>> gtext('alfa2')
>> gtext('alfa1')
>> gtext('mye=0.2')

```

The result is presented in the fig.3

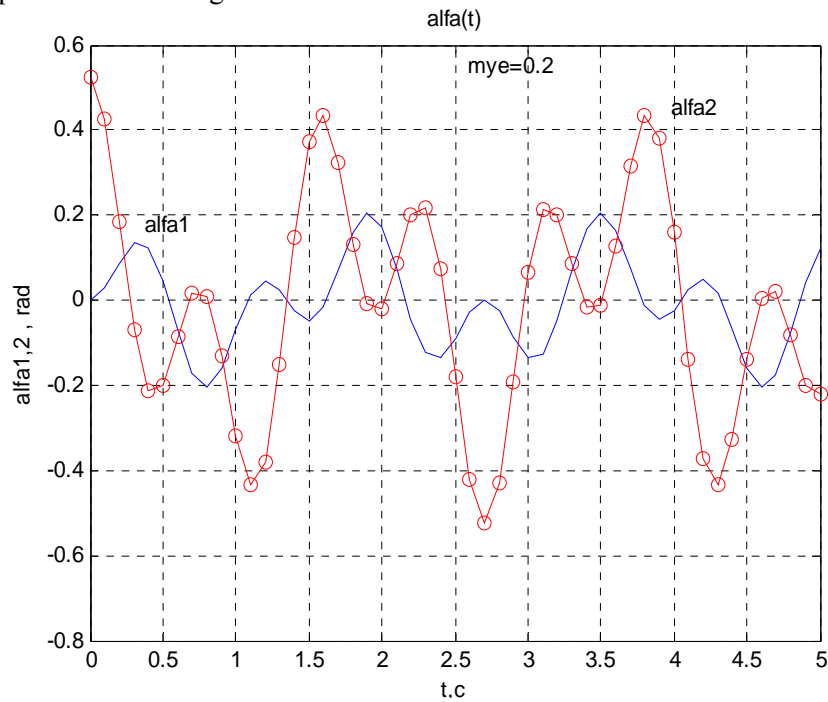


Fig.3 - The graphs of small oscillations of the double pendulum at  $\mu_2 = 0.2$

```

>> mye=0.3;
>> g=9.8;
>> l=0.25;
>> A=mye+((1+mye)*mye).^0.5;
>> w1=(g/l).^0.5*(1+A).^0.5;
>> w2=(g/l).^0.5*(1-A).^0.5;
>> t=0:0.1:5;
>> alfa1=(-pi/12)*(mye./(1+mye)).^0.5*(cos(w1.*t))+...
(pi/12)*(mye./(1+mye)).^0.5*(cos(w2.*t));
>> plot(t,alfa1)
>> grid on
>> hold on

```

```

>>alfa2=(pi/12)*cos(w1.*t)+(pi/12)*cos(w2.*t);
>>plot(t,alfa2,'ro-')
>>gtext('alfa2')
>> gtext('alfa1')
>> gtext('mye=0.3')

```

The result is presented in the fig.4

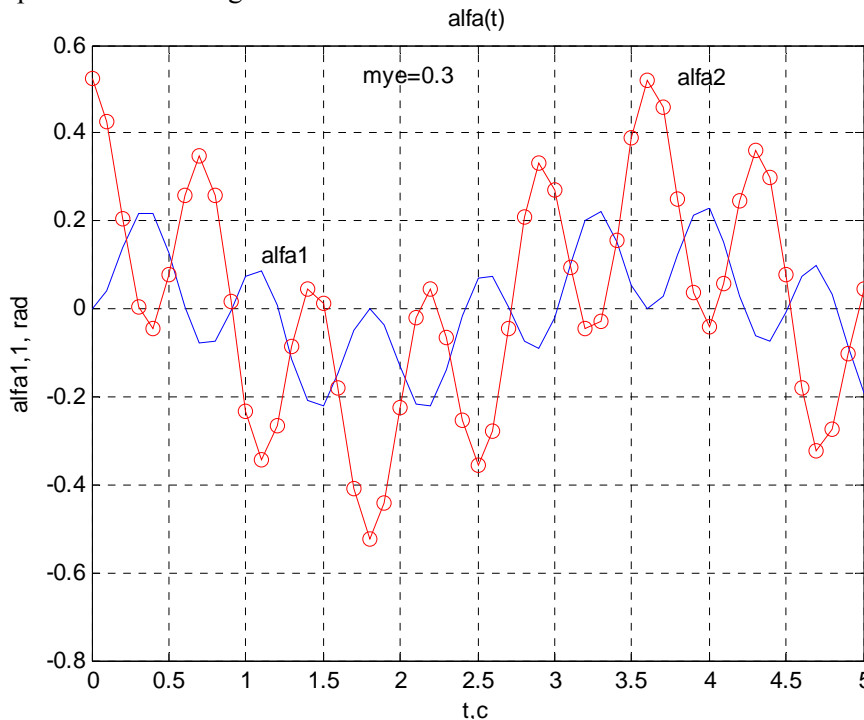


Fig.4 - The graphs of small oscillations of the double pendulum at  $\mu_3 = 0.3$

Here angles  $\alpha_1(t)$ ,  $\alpha_2(t)$  are expressed in radians, and  $t$  time is taken in seconds. Figures 2-4 demonstrate the graphs of small oscillations of pendulums for three values of  $\mu$ :  $\mu_1=0.1$ ,  $\mu_2=0.2$ ,  $\mu_3=0.3$  when  $l_1 = l_2 = l = 0.25m$ ,  $g = 9.8 m/s^2$ . The angles of displacement of pendulums from equilibrium position are given in radians. It is seen from the graphs that beats occur during double pendulum small oscillations at which energy is cyclically transmitted from one pendulum to the other one. When one pendulum almost stops, the other one is at its maximum amplitude. After a while pendulums "change roles" and so on. Oscillations with a bigger frequency  $\omega_1$  are modulated by lower-frequency oscillations with frequency  $\omega_2$ .

**Conclusion.** The article considers the calculation and visualization of small oscillations of a double plane simple pendulum. It contains the brief derivation of the motion equation and its solutions; the mathematical model of the motion in the form of the system of nonlinear differential equations. Experiments with the double pendulum are performed at various ratios of bodies' masses and initial angles, in particular for three values of  $\mu = m_1 / m_2$ :  $\mu_1= 0.1$ ,  $\mu_2= 0.2$ ,  $\mu_3= 0.3$  when the lengths of the pendulum are  $l = l_1 = l_2 = 0.25m$  and  $g = 9.8 m/s^2$ . The angles of pendulum deviation are given in radians. The graphs of small oscillations of the pendulum show that beats occur in the system during which energy cyclically passes from one pendulum to another. When one pendulum almost stops, the other pendulum is at its maximum amplitude. After a while pendulums "exchange their states" and so on. Oscillations with a bigger frequency  $\omega_1$  are modulated by lower frequency oscillations with a frequency  $\omega_2$ . The results of calculation and visualization of small oscillations of a double plane simple pendulum can be used in the theoretical mechanics.

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## ҚОС ЖАЗЫҚ МАЯТНИКТИҢ АУЫТҚУЫ КІШІ ТЕРБЕЛІСТЕРІН ЕСЕПТЕУ МЕН БЕЙНЕЛЕУ

**Аннотация.** Қос жазық математикалық маятниктің ауытқуы кіші тербелістерін есептеу мен бейнелеу ұсынылады. Қозғалыс теңдеулерін қорытып шығарылуы мен шешімі келтірілген, сызықтық емес дифференциал теңдеулер бойынша математикалық моделі құрастырылған. Денелердің массаларының әртүрлі қатынастары мен бастапқы ауытқу бұрыштары әртүрлі, атап айтқанда,  $g = 9.8\text{м/с}^2$ ,  $l=l_1=l_2=0.25\text{м}$  шартында, массалар қатынасының үш түрі шамасында  $\mu = m_1 / m_2$ :  $\mu_1=0.1$ ,  $\mu_2=0.2$ ,  $\mu_3=0.3$  жағдайлары үшін қос маятниктің қозғалысын бақылауға арналған эксперименттер жүргізілген. Маятниктердің кіші тербелістерінің графиктері келтірілген. Ауытқу бұрыштары радианда берілген.

Графиктерден жүйеде маятниктердің арасында циклді энергия алмасуы нәтижесінде соққы құбылысы пайда болатыны байқалады. Маятниктің біреуінің ауытқуы тоқтағанда екіншісі максимал амплитудада тербеле бастайды. Әлдебір уақыттан кейін маятниктер рөлі ауысады және осылайша қайталанатын. Тербеліс жиілігі  $\omega_1$  үлкені жиілігі кіші  $\omega_2$  тербеліспен модуляцияланады.

**Түйін сөздер.** Қос маятник, ауытқуы кіші тербелістер, соққы, энергиямен алмасу, меншікті жиіліктер, нормаль (табиғи) модалар.

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## РАСЧЕТ И ВИЗУАЛИЗАЦИЯ МАЛЫХ КОЛЕБАНИЙ ДВОЙНОГО ПЛОСКОГО МАЯТНИКА

**Аннотация.** Предлагается расчет и визуализация малых колебаний двойного плоского математического маятника. Приведен краткий вывод уравнения движения и их решения, построена математическая модель в виде системы нелинейных дифференциальных уравнений. Проведены эксперименты по наблюдению за движением двойного маятника при различных отношениях масс тел и начальных углах, в частности для трех значений  $\mu = m_1 / m_2$ :  $\mu_1=0.1$ ,  $\mu_2=0.2$ ,  $\mu_3=0.3$ , при условии  $l=l_1=l_2=0.25\text{м}$ ,  $g = 9.8\text{м/с}^2$ . Приведены графики малых колебаний маятников. Углы отклонения маятников приведены в радианах. Из графиков видно, что в системе происходят биения, при которых энергия циклически переходит от одного маятника к другому. Когда один маятник почти останавливается, другой раскачивается с максимальной амплитудой. Через некоторое время маятники "меняются ролями" и так далее. Колебания с большей частотой  $\omega_1$  модулируются низкочастотными колебаниями с частотой  $\omega_2$ .

**Ключевые слова.** Двойной маятник, малые колебания, биения, обмен энергиями, характерные частоты, нормальные мода.

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**ISSN 2518-1726 (Online), ISSN 1991-346X (Print)**

Редакторы *М. С. Ахметова, Т.А. Апендиев, Д.С. Аленов*  
Верстка на компьютере *А.М. Кульгинбаевой*

Подписано в печать 10.04.2019.  
Формат 60x881/8. Бумага офсетная. Печать – ризограф.  
5,8 п.л. Тираж 300. Заказ 2.

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*Национальная академия наук РК*  
*050010, Алматы, ул. Шевченко, 28, т. 272-13-18, 272-13-19*